

A Modification and Analysis of Lagrangian Trajectory Modeling and Granular Dynamics of Lunar Dust Particles

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Abstract

A previously developed mathematical model is amended to more accurately incorporate the effects of lift and drag on single dust particles in order to predict their behavior in the wake of high velocity gas flow. The model utilizes output from a CFD or DSMC simulation of exhaust from a rocket nozzle hot gas jet. An extension of the Saffman equation for lift based on the research of McLaughlin (1991) and Mei (1992) is used, while an equation for the Magnus force modeled after the work of Oesterle (1994) and Tsuji et al (1985) is applied. A relationship for drag utilizing a particle shape factor ($\phi = 0.8$) is taken from the work of Haider and Levenspiel (1989) for application to non-spherical particle dynamics. The drag equation is further adjusted to account for rarefaction and compressibility effects in rarefied and high Mach number flows according to the work of Davies (1945) and Loth (2007) respectively. Simulations using a more accurate model with the correction factor ($C = 0.8$ in a 20% particle concentration gas flow) given by Richardson and Zaki (1954) and Rowe (1961) show that particles have lower ejection angles than those that were previously calculated. This is more prevalent in smaller particles, which are shown through velocity and trajectory comparison to be more influenced by the flow of the surrounding gas. It is shown that particles are more affected by minor changes to drag forces than larger adjustments to lift forces, demanding a closer analysis of the shape and behavior of lunar dust particles and the composition of the surrounding gas flow.

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The diagram illustrates a simulation workflow. At the top, a box labeled 'CFD (Computation Fluid Dynamics) or DSMC (direct simulation Monte Carlo)' is connected by a blue arrow to a box labeled 'CFD: Sophisticated COTS Software to Solve Navier Stokes Equations, which describes motion of a gas or liquid'. This box is then connected by a red arrow to a box labeled 'Particle Ballistics Simulation'. From the 'Particle Ballistics Simulation' box, two arrows emerge: a blue arrow pointing to a box labeled 'Gas:' and a red arrow pointing to a box labeled 'Particle:'. The 'Gas:' box contains a list of properties: Density, Velocity, and Temperature. The 'Particle:' box contains a list of properties: Forces, Acceleration, Velocity, and Position. Below the 'Particle Ballistics Simulation' box, the text 'Generates Particle Trajectories' is followed by 'Particle' and a list of parameters: Diameter and Initial Starting Point.

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graph TD
    A[CFD (Computation Fluid Dynamics) or DSMC (direct simulation Monte Carlo)] -- blue arrow --> B[CFD: Sophisticated COTS Software to Solve Navier Stokes Equations, which describes motion of a gas or liquid]
    B -- red arrow --> C[Particle Ballistics Simulation]
    C -- blue arrow --> D[Gas:  
• Density  
• Velocity  
• Temperature]
    C -- red arrow --> E[Particle:  
• Forces  
• Acceleration  
• Velocity  
• Position]
    C --> F[Generates Particle Trajectories  
Particle  
• Diameter  
• Initial Starting Point]
  
```

Key Point #2

Assumption:

Ballistics calculations ignore the temporal variation of the gas.

Time Scale of Lander Descent >> Time Scale of Particle Trajectory

This is analogous to ignoring the effect of grass growing along the path of a golf ball rolling across a putting green

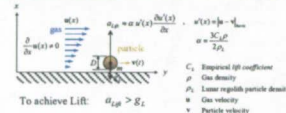

$$\begin{aligned}v_n &= v_{n-1} + \Delta t a_{n-1} \\r_n &= r_{n-1} + \Delta t v_{n-1} + \frac{1}{2} \Delta t^2 a_{n-1} \\a_n &= F(r_n, v_n) / m\end{aligned}$$

Simple Model

Vertical Gradient of Horizontal Force $\neq 0$
 \Rightarrow Bernoulli Lift Force

If any one of these 3 quantities go to 0, the *Bernoulli Lift Force* also goes to 0

- *Relative gas velocity*
- *Vertical gradient of relative gas velocity*
- *Gas density*

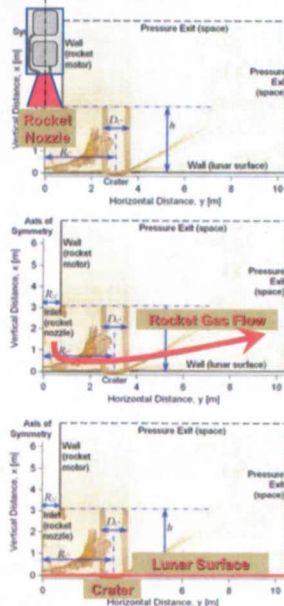
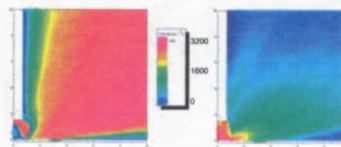
[illegible]

Using the method described above, integrated along a sphere's horizontal axis and arrives at an expression that includes two terms:

$$F_L = 8\pi\eta b^3(u-v)\left(\omega_b + \frac{1}{2}\frac{d\omega}{dv}\right)$$

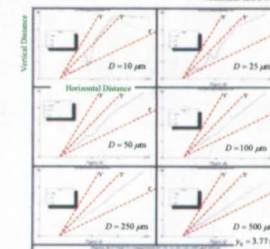
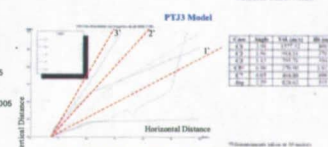
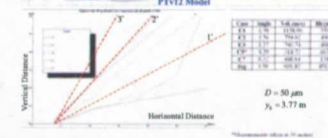
where,
 u = the gas velocity
 v = particle velocity
 ρ = gas density
 b = particle radius
 ω_s = sphere's angular rotation speed
 du/dy = vertical gradient of horizontal component of gas speed

- The 1st term looks like the Magnus force, the 2nd term looks like the Saffman force
- If $u_0 = -1/2 \, \text{sh/sy}$, then there is no lift, as would be expected.

CFD Gas Velocity [m/s] CFD Gas Density [kg/m³]

Height of Nozzle above Surface, h	Distance from Nozzle Center Line to Crater Center, R_c	
	10 ft	30 ft
5 ft	Case C2	---
10 ft	Case C7	Case C1
20 ft	Case C1	---

- **PT12 - Original / Lane and Metzger (Mar 2000)**
- Lane and Metzger (Jun - Aug 2006)
- **PT14 - Drug / Modifications:** Hailer & Levenapfel Model with Richardson & Zaki Correction Factor. This version of the code modifies drug force equations only
- **PT15 - LbH Modifications:** Extended Saffman LbH Force. This version of the code modifies LbH equations only, and the extended Saffman LbH equation.
- **PT19 - Drug / LbH Modifications:** This version merges PT14 and PT15.
- **PT16 - Drug / LbH Modifications:** Compressibility/Extended Effects with Hailer & Levenapfel Model, Magnus Effect, and Extended Saffman LbH Force.



Point #1

Ballistic Angles of simulated particle trajectories \approx Trajectory Angles derived from Apollo video analysis

- Craters can generate a substantial disruption to particle ballistics motion.
- Particles can easily fall into craters.
- Particle motion down-stream from craters, at the surface layer, may be suppressed.

Comparison of Integration Algorithms

- **2nd Order Taylor Series** algorithm is very stable and well suited for single particle simulations, as long as the step-size is sufficiently small.
- **Beeman's Method** is Much Better
- **RK4** is superior to both

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